Linear Systems, 2019 - Lecture 2

- Transition Matrix Properties
- Time-varying change of coordinates
- Periodic Systems
- Floquet Decomposition
- Time-varying Transfer Functions

Rugh, Chapter 5 [and Chapter 21]

Main news:

- Properties of LTV systems
- LTP systems

Continuous Time-varying (LTV) Systems

For bounded A(t), the equation

$$\dot{x}(t) = A(t)x(t), \quad x(t_0) = x_0$$

has a unique solution of the form

$$x(t) = \Phi(t, t_0) x_0$$

The *transition matrix* $\Phi(t, t_0)$ can be written as the infinite sum

$$\Phi(t, t_0) = I + \int_{t_0}^t A(\sigma_1) d\sigma_1 + \int_{t_0}^t A(\sigma_1) \int_{t_0}^{\sigma_1} A(\sigma_2) d\sigma_2 d\sigma_1 + \int_{t_0}^t A(\sigma_1) \int_{t_0}^{\sigma_1} A(\sigma_2) \int_{t_0}^{\sigma_2} A(\sigma_3) d\sigma_3 d\sigma_2 d\sigma_1$$

Transition Matrix $\Phi(t, t_0)$

The unique solution of the equation

$$\frac{d}{dt}X(t) = A(t)X(t)$$
$$X(t_0) = I$$

is $X(t) = \Phi(t, t_0)$.

Proof. Let $x(t) = X(t)x_0$. Then

$$\dot{x}(t) = \frac{d}{dt}X(t)x_0 = A(t)X(t)x_0 = A(t)x(t)$$

so

$$x(t) = \Phi(t, t_0) x_0$$

Hence $\Phi(t,t_0)x_0 = X(t)x_0$ for every x_0 , so $\Phi(t,t_0) = X(t)$

Nice Example: Scalar Time-variation

Consider

$$\dot{x} = Aa(t)x(t)$$

The transition matrix is

$$\Phi(t,t_0) = I + A \int_{t_0}^t a(\sigma_1) d\sigma_1 + A^2 \int_{t_0}^t a(\sigma_1) \int_{t_0}^{\sigma_1} a(\sigma_2) d\sigma_2 d\sigma_1 + \cdots$$
$$= \sum_{k=0}^\infty \frac{1}{k!} A^k \left[\int_{t_0}^t a(\sigma) d\sigma \right]^k$$
$$= \exp\left(A \int_{t_0}^t a(\sigma) d\sigma\right)$$

Second equality is nontrivial.

(Recall Two Tank Example with time-varying flow q(t))

More general case: Commutating A(t)

lf

$$A(t)\int_{t_0}^t A(\sigma)d\sigma = \int_{t_0}^t A(\sigma)d\sigma A(t)$$

then

$$\Phi(t,t_0) = \exp\left\{\int_{t_0}^t A(\sigma)d\sigma\right\}$$

Special case: $A(t)A(\tau) = A(\tau)A(t)$ for all t,τ

Example

If $A(t) = a_1(t)A_1 + a_2(t)A_2$ where A_1 and A_2 commute then

$$\begin{split} \Phi(t,t_0) &= & \exp\left\{\int_{t_0}^t a_1(t)A_1 + a_2(t)A_2dt\right\} \\ &= & \exp\left\{\int_{t_0}^t a_1(t)dtA_1\right\}\exp\left\{\int_{t_0}^t a_2(t)dtA_2\right\} \end{split}$$

Example

$$\begin{aligned} \dot{x}(t) &= \begin{bmatrix} 1 & \cos t \\ 0 & 0 \end{bmatrix} x(t) \\ x_2(t) &\equiv x_2(\tau) \\ \dot{x}_1(t) &= x_1(t) + \cos t \cdot x_2(\tau) \\ x_1(t) &= e^{t-\tau} x_1(\tau) + \int_{\tau}^{t} e^{t-\sigma} \cos \sigma d\sigma \cdot x_2(\tau) \\ &= e^{t-\tau} x_1(\tau) + \frac{1}{2} \left(\sin t - \cos t - e^{t-\tau} (\sin \tau - \cos \tau) \right) \cdot x_2(\tau) \\ \Phi(t,\tau) &= \begin{bmatrix} e^{t-\tau} & \frac{1}{2} \left(\sin t - \cos t - e^{t-\tau} (\sin \tau - \cos \tau) \right) \\ 0 & 1 \end{bmatrix} \end{aligned}$$

Sanity check:
$$\Phi(t,t) = I$$
 and $\frac{d}{dt}\Phi(t,\tau)\Big|_{t=\tau} = \begin{bmatrix} 1 & \cos t \\ 0 & 0 \end{bmatrix}$

7/41

The equation

$$\begin{aligned} \dot{x}(t) &= A(t)x(t) + B(t)u(t) \\ x(t_0) &= x_0 \end{aligned}$$

has the unique solution

$$x(t) = \Phi(t, t_0)x_0 + \int_{t_0}^t \Phi(t, \sigma)B(\sigma)u(\sigma)d\sigma$$

Proof: Differentiate!

Properties of $\Phi(t,\sigma)$

For any t, τ, σ , the transition matrix satisfies

$$\Phi(t,\tau) = \Phi(t,\sigma)\Phi(\sigma,\tau) \text{ (semigroup property)}$$
$$\frac{d}{dt}\Phi(t,\sigma) = A(t)\Phi(t,\sigma)$$
$$\frac{d}{d\sigma}\Phi(t,\sigma) = -\Phi(t,\sigma)A(\sigma)$$

Proof of first property: Let $R(t) = \Phi(t, \sigma) \Phi(\sigma, \tau)$. Then

$$\frac{d}{dt}R(t) = A(t)R(t)$$
$$R(\sigma) = \Phi(\sigma, \tau)$$

so R(t) must be identical to $\Phi(t,\tau)$

Properties of $\Phi(t,\sigma)$

Proof of third property:

$$\Phi(\sigma + h, \sigma) = I + hA(\sigma) + o(h) \qquad \text{(why?)}$$

Hence, using first property, we have

$$\Phi(t,\sigma) = \Phi(t,\sigma+h)(I+hA(\sigma)+o(h))$$

from which we get

$$\frac{1}{h}(\Phi(t,\sigma+h) - \Phi(t,\sigma)) = -\Phi(t,\sigma+h)A(\sigma) + o(1)$$

from which the result follows as $h \to 0$

$$\frac{d}{d\sigma}\Phi(t,\sigma) = -\Phi(t,\sigma)A(\sigma)$$

Inversion

The transition matrix $\Phi(t,t_0)$ is invertible for any t,t_0 and

$$\Phi(t, t_0)^{-1} = \Phi(t_0, t)$$

Proof. By the composition rule

$$\Phi(t, t_0)\Phi(t_0, t) = \Phi(t_0, t)\Phi(t, t_0) = \Phi(t_0, t_0) = I$$

For any $t,\tau,\sigma,$ the transition matrix satisfies

$$\Phi(t,t) = I$$

$$\Phi(t,\tau) = \Phi(t,\sigma)\Phi(\sigma,\tau)$$

$$\Phi(t,\sigma))^{-1} = \Phi(\sigma,t)$$

$$\frac{d}{dt}\Phi(t,\sigma) = A(t)\Phi(t,\sigma)$$

$$\frac{d}{d\sigma}\Phi(t,\sigma) = -\Phi(t,\sigma)A(\sigma)$$

Change of Variables

Variable change x(t) = P(t)z(t) (with P(t) invertible) gives

$$\dot{x}(t) = A(t)x(t), \quad x(t_0) = x_0 \iff \dot{z}(t) = \left[P(t)^{-1}A(t)P(t) - P(t)^{-1}\dot{P}(t) \right] z(t), \quad z(t_0) = P(t)^{-1}x_0$$

For the fundamental matrix this means that

$$\Phi_{P^{-1}AP - P^{-1}\dot{P}}(t, t_0) = P^{-1}(t)\Phi_A(t, t_0)P(t_0)$$

Proof:

$$\begin{aligned} APz &= Ax &= \dot{x} = \dot{P}z + P\dot{z} \\ z(t) &= \Phi_{P^{-1}AP - P^{-1}\dot{P}}(t, t_0)z(t_0) \\ z(t) &= P^{-1}(t)x(t) = P^{-1}(t)\Phi_A(t, \tau)P(t_0)z(t_0) \end{aligned}$$

Adjoint system

From
$$\frac{d}{d\sigma}\Phi_A(t,\sigma) = -\Phi_A(t,\sigma)A(\sigma)$$
 follows that $\Phi_A^T(t,\sigma)$ solves
$$\frac{d}{d\sigma}Z(\sigma) = -A^T(\sigma)Z(\sigma), \quad Z(t) = I$$

This fact can be written as

$$\Phi_{-A^T}(t,t_0) = \Phi_A^T(t_0,t)$$

Define $\boldsymbol{X}(\boldsymbol{k})$ recursively as

$$\begin{aligned} X(k+1) &= A(k)X(k), \quad k \geq k_0 \\ X(k_0) &= I \end{aligned}$$

Then $\Phi(k, k_0) = X(k)$.

Remark: What about $\Phi(k, k_0)$ when $k < k_0$?

The example $x(k+1) = 0 \cdot x(k)$ shows that x(k) might not be uniquely determined by $x(k_0)$ for $k < k_0!$

Difference between discrete and continuous time

Properties of $\Phi(k, k_0)$

$$\begin{split} \Phi(k+1,j) &= A(k)\Phi(k,j), \quad k \ge j \\ \Phi(k,j-1) &= \Phi(k,j)A(j-1), \quad k \ge j \\ \Phi(k,i) &= \Phi(k,j)\Phi(j,i), \quad k \ge j \ge i \end{split}$$

If the $n \times n$ matrix A(k) is invertible for each k, then $\Phi(k, j)$ is invertible for each $k \ge j$ and $\Phi(j, k)$ can be defined as

$$\Phi(j,k) = \Phi(k,j)^{-1}$$

Change of Variables

Variable change x(k) = P(k)z(k) (with P(k) invertible) we get

$$\begin{aligned} x(k+1) &= A(k)x(k), \quad x(k_0) = x_0 &\iff \\ z(k+1) &= \left[P(k+1)^{-1}A(k)P(k) \right] z(k) \end{aligned}$$

Hence we have

$$\Phi_z(k,j) = P(k)^{-1} \Phi_x(k,j) P(j)$$

Theorem by Abel-Jacobi-Liouville

Let A(t) be continuous. Then

$$\det \Phi(t, t_0) = \exp\left(\int_{t_0}^t \operatorname{tr}[A(\sigma)] d\sigma\right)$$

Interpretation: Volume contraction

Proof: Let c_{ij} be the cofactor of entry ϕ_{ij}

$$\frac{d}{dt}\det\Phi(t,t_0) = \sum_{i,j} \left(\frac{\partial}{\partial\phi_{ij}}\det\Phi(t,t_0)\right)\dot{\phi}_{ij}(t,t_0)$$

$$= \sum_{i,j} c_{ij}(t,t_0)\dot{\phi}_{ij}(t,t_0)$$

$$= \operatorname{tr}\left(C(t,t_0)^T\dot{\Phi}(t,t_0)\right)$$

$$= \operatorname{tr}\left(\Phi(t,t_0)C(t,t_0)^TA(t)\right)$$

$$= \operatorname{tr}\left((\det\Phi(t,t_0)I)A(t)\right)$$

$$= \operatorname{tr}A(t) \cdot \det\Phi(t,t_0)$$

Is it possible to asymptotically stabilize the oscillative system

$$\dot{x}(t) = \begin{bmatrix} 0 & 1 \\ -1 & 0 \end{bmatrix} x + \begin{bmatrix} 0 \\ 1 \end{bmatrix} u$$
$$y = \begin{bmatrix} 1 & 0 \end{bmatrix} x$$

by time-varying output feedback

$$u(t) = -l(t)y(t)?$$

Solution

The closed loop system becomes $\dot{x} = A_c x$ with

$$A_c = \begin{bmatrix} 0 & 1\\ -1 - l(t) & 0 \end{bmatrix}$$

By the Abel-Liouville theorem we have

$$\det \Phi(t,0) = \exp(t \operatorname{tr} A_c) \equiv 1$$

The system can hence not be asymptotically stable, since an asymptotically stable system must have $\Phi(t,0) \to 0$ as $t \to \infty$ [Why?].

Consider the time-varying system

$$\dot{x} = e^{-At} B e^{At} x \tag{1}$$

Note that $e^{-At}Be^{At}$ has the same eigenvalues as B

The coordinate change $\boldsymbol{z}(t)=e^{At}\boldsymbol{x}(t)$ transforms the system to

$$\dot{z} = (A+B)z\tag{2}$$

Example - LTV Systems and Eigenvalues

Assume

$$A = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix}$$

Then with $z(t) = e^{At}x(t)$ one has ||z|| = ||x||, so asymptotic stability of (1) and (2) are equivalent (rotating coordinate transformation).

Proof: More generally if A(t) is skew-symmetric for all t, i.e. $A^T + A = 0$ then $Q(t) = \Phi_A(t, 0)$ is orthogonal, i.e. satisfies $Q^T Q = I$, since

$$\frac{d}{dt}Q^TQ = Q^T(A^T + A)Q = 0, \qquad Q(0) = I$$

Therefore $z^T z = x^T Q^T Q x = x^T x$.

Example - LTV Systems and Eigenvalues

With the stable matrix

$$B = \begin{bmatrix} -1 & M \\ 0 & -1 \end{bmatrix}$$

it is easy to see that A + B is unstable for M > 2.

Hence system (1) above is an **unstable time-varying system with stable eigenvalues** (equal to -1 for all t).

Example -LTV Systems and Eigenvalues

With the unstable matrix

$$B = \begin{bmatrix} -1 & 0\\ 0 & 1/2 \end{bmatrix}$$

it is easy to see that A + B is stable.

Hence system (1) above is a stable LTV system having one unstable eigenvalue for all t.

Exercise: Can you find a 2×2 asymptotically stable LTV system with both eigenvalues in the RHPL for all t?

Linear Time Periodic (LTP) Systems

A linear system

$$\dot{x}(t) = A(t)x(t)$$

with

$$A(t+T) = A(t)$$

is said to be T-periodic.

The smallest such T is called the period of the system.

A state space system is called T-periodic if all matrices $\left(A,B,C,D\right)$ are T-periodic.

The following is the main result for periodic systems

Floquet Decomposition

"Long-term trend + periodic fluctuations"

Let A(t) be bounded and T-periodic. Then for

 $\dot{x}(t) = A(t)x(t), \quad x(t_0) = x_0$

the transition matrix can be written

$$\Phi(t,\tau) = P(t)e^{R(t-\tau)}P(\tau)^{-1}$$

where $R \in \mathbf{C}^{n \times n}$ is constant and $P(t) \in \mathbf{C}^{n \times n}$ is differentiable, invertible, and T-periodic.

The variable transformation x(t) = P(t)z(t) gives $\dot{z} = Rz$

Proof

Since $\Phi(T,0)$ is nonsingular, there exists a solution $F \in C^{n \times n}$ (in fact infinitly many) to $e^F = \Phi(T,0)$. Choosing any such F, define $R = \frac{1}{T}F$, we then have

$$e^{RT} = \Phi(T,0)$$

Define then P(t) by

$$P(t) = \Phi(t,0)e^{-Rt}$$

We get

$$\begin{split} \Phi(t,\tau) &= \Phi(t,0)\Phi(\tau,0)^{-1} = P(t)e^{R(t-\tau)}P(\tau)^{-1} \\ P(t+T) &= \Phi(t+T,0)e^{-R(t+T)} \\ &= \Phi(t+T,T)\Phi(T,0)e^{-RT}e^{-Rt} \\ &= \Phi(t+T,T)e^{-Rt} \\ &= \Phi(t,0)e^{-Rt} \\ &= P(t) \end{split}$$

Discrete Time Floquet Decomposition

Let ${\cal A}(k)$ be $K\mbox{-periodic}.$ Then for

$$x(k+1) = A(k)x(k), \quad x(k_0) = x_0$$

the transition matrix can be written

$$\Phi(k,j) = P(k)R^{(k-j)}P(j)^{-1}$$

where $R \in \mathbb{C}^{n \times n}$ and P(k) is *K*-periodic.

With x(k) = P(k)z(k), this gives

$$z(k+1) = Rz(k)$$

2-periodic Example

$$A(k) = \begin{bmatrix} (-1)^{k} & 0\\ 0 & 1 \end{bmatrix}$$
$$R^{2} = \Phi(2,0) = \begin{bmatrix} -1 & 0\\ 0 & 1 \end{bmatrix}$$
$$R = \begin{bmatrix} i & 0\\ 0 & 1 \end{bmatrix}$$

Note that R is not real-valued !

Extra: Real Floquet Factors for LTP Systems

(not in Rugh)

It is always possible to obtain a **real** Floquet factorisation for a real T-periodic system, by treating the system as having 2T-periodic coefficients:

From the fact that $\Phi(2T,0) = \Phi(T,0)^2$ it can be proved (use Jordan-form) that there is a real matrix G such that

$$e^{2TG} = \Phi(2T, 0).$$

Then $P(t) := \Phi(t, 0)e^{-tG}$ is real and can as before be seen to be 2T-periodic (but not necessarily T-periodic).

See Montagnier, P, et.al Real Floquet Factors of Linear Time-Periodic Systems (Google it)

LTI with sinusodal input - Resonances

Consider the equation

$$\dot{x} = \begin{bmatrix} 0 & -1 \\ 1 & 0 \end{bmatrix} x + \begin{bmatrix} 0 \\ \sin t \end{bmatrix} \quad x(0) = x_0$$

Laplace transform:

$$\mathbf{x}_{2}(s) = C(sI - A)^{-1}(Bu(s) + x_{0})$$

= $\frac{s}{(1 + s^{2})^{2}} + \frac{1}{1 + s^{2}} \begin{bmatrix} 1 & s \end{bmatrix} x_{0}$

$$x_2(t) = \frac{t}{2} \sin t + \begin{bmatrix} \sin t & \cos t \end{bmatrix} x_0$$

For what systems does periodic input give periodic solution?

For $A \in \mathbf{R}^{n \times n}$ and

$$\dot{x}(t) = Ax(t) + f(t)$$

one can prove that the following statements are equivalent:

- (i) No eigenvalue of A has zero real part.
- (ii) A unique T-periodic solution exists for every T-periodic f.

Theorem 5.15

Let A(t) be continuous and T-periodic and

$$\dot{x}(t) = A(t)x(t) + f(t)$$

The following statements are then equivalent:

- (i) No nontrivial *T*-periodic solution exists for $f \equiv 0$.
- (ii) A unique T-periodic solution exists for every T-periodic f.

Time-varying Transfer Functions

(not in Rugh)

Transfer function analysis is quite involved for time-varying linear systems

$$y(t) = \int h(t,\tau) u(\tau) d\,\tau$$

For LTI systems, $h(t,\tau)=h(t-\tau,0)$ and

$$Y(j\omega) = H(j\omega)U(j\omega), \qquad H(j\omega) = \int h(r,0)e^{-j\omega r}dr$$

What can be said for general LTV systems?

And for LTP systems, where $h(t+T,\tau+T) = h(t,\tau)$?

Transfer Functions for LTV Systems

Define as usual

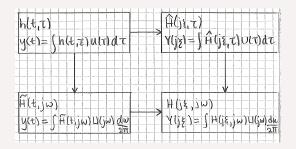
$$U(j\omega) = \int e^{-j\omega\tau} u(\tau) d\tau; \qquad u(t) = \int e^{j\omega\tau} U(j\omega) d\omega/2\pi$$
$$Y(j\xi) = \int e^{-j\xi t} y(t) dt; \qquad y(t) = \int e^{j\xi t} Y(j\xi) d\xi/2\pi$$

Define also

$$\begin{aligned} \widehat{H}(j\xi,\tau) &= \int e^{-j\xi t} h(t,\tau) dt \\ \widetilde{H}(t,j\omega) &= \int e^{j\omega\tau} h(t,\tau) d\tau \\ H(j\xi,j\omega) &= \iint e^{j\omega\tau-j\xi t} h(t,\tau) d\tau dt \end{aligned}$$

We then have the following relations (if convergence ok)

Transfer Functions for LTV Systems



For LTI systems we have

$$H(j\xi, j\omega) = 2\pi\delta(\xi - \omega)H(j\omega)$$

and

$$Y(j\omega) = H(j\omega)U(j\omega)$$

Transfer Functions for LTP Systems

If
$$h(t+T,\tau+T) = h(t,\tau)$$
 then (with $\omega_0 = 2\pi/T$)

$$Y(j\omega) = \sum_{k} H_k(j\omega - jk\omega_0)U(j\omega - jk\omega_0)$$

where

$$H_k(j\omega) = F(h_k(t)) = \int e^{-j\omega t} h_k(t) dt$$

with

$$h_k(t) = \frac{1}{T} \int_0^T h(r, r-t) e^{-jk\omega_0 r} dr$$

The interpretation is that H_k describes the signal transfer from input frequency ω to output frequency $\omega + k\omega_0$.

Next Week

- Controllability and Observability
- Controller/Observer Forms
- Balanced Realizations

Bonus: Abel-Jacobi-Liouville alt. proof

Want to prove that

$$\frac{d}{dt}\det\Phi(t,t_0) = \operatorname{tr} A(t)\det\Phi(t,t_0)$$

Taylor-expansion gives

$$\Phi(t + dt, t_0) = \Phi(t, t_0) + A(t)\Phi(t, t_0)dt + o(dt)$$

= $(I + A(t)dt)\Phi(t, t_0) + o(dt)$

Since the determinant is the product of the eigenvalues, and these satisfy $\lambda_i(I+Adt)=1+\lambda_i(A)dt$ we get

$$\det\Phi(t+dt,t_0) = \prod(1+\lambda_i(A)dt) \cdot \det\Phi(t,t_0) + o(dt)$$
$$= (1+\sum\lambda_i(A)dt + o(dt)) \cdot \det\Phi(t,t_0) + o(dt)$$

From which the result follows by letting $dt \rightarrow 0$.

The unstable equilibrium becomes stable if the point of suspension oscillates fast in the vertical direction:

Pendulum length l

Oscillation amplitude a << l

Period of oscillation 2τ

Acceleration supposed *constant* equal to $\pm c$ (so $c = 8a/\tau^2$)

Assume

$$\ddot{x} = (\omega^2 \pm d^2)x$$

where the sign changes after time τ , where $\omega^2 = g/l$ and $d^2 = c/l$. If the oscillation of the suspension is fast enough, c > g, then $d^2 = 8a/(l\tau^2) > \omega^2$.

Bonus: The Oscillating Inverted Pendulum

If you have time over:

[a.] Show that $\Phi(2\tau,0)=A_2A_1$

$$A_1 = \begin{bmatrix} \cosh k\tau & 1/k \sinh k\tau \\ k \sinh k\tau & \cosh k\tau \end{bmatrix}, \quad A_2 = \begin{bmatrix} \cos\Omega\tau & 1/\Omega \sin\Omega\tau \\ -\Omega \sin\Omega\tau & \cos\Omega\tau \end{bmatrix}$$

where
$$k^2 = d^2 + \omega^2$$
 and $\Omega^2 = d^2 - \omega^2$.

[b.] Show that $\det(\Phi(t,0)) = 1$, $t = 2\tau, 4\tau, 6\tau, \ldots$ and that the criterion for stability is $|tr(\Phi(2\tau,0))| < 2$

[c.] Suppose that $a \ll l$. Show that the system is stable for sufficiently large $c \gg g$. The stability condition is approximately

 $c/g > (3/2) \, l/a$